

# A "Gauss–Bonnet Formula" for Contact Sub-Riemannian Manifolds

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## Abstract

We study 3-dimensional manifolds endowed with oriented contact sub-Riemannian structures. The Euler characteristic class of the contact structure is presented as the *rotation class* of a volume preserving vector field constructed in terms of fundamental differential invariants of the sub-Riemannian metric.

**1.** Let  $M$  be a smooth 3-dimensional manifold. A *contact sub-Riemannian structure* is a pair  $\Delta, \langle \cdot | \cdot \rangle$ , where  $\Delta = \{\Delta_q\}_{q \in M}$ ,  $\Delta_q \subset T_q M$ , is a contact structure on  $M$  and  $\langle \cdot | \cdot \rangle = \{\langle \cdot | \cdot \rangle_q\}_{q \in M}$  is a smooth with respect to  $q$  family of Euclidean inner products

$$(v_1, v_2) \mapsto \langle v_1 | v_2 \rangle_q, \quad v_1, v_2 \in \Delta_q,$$

defined on  $\Delta_q$ . A Lipschitzian curve  $\xi : [0, 1] \rightarrow M$  is called *admissible* for  $\Delta$  if  $\frac{d\xi(t)}{dt} \in \Delta_{\xi(t)}$  for almost all  $t \in [0, 1]$ . The *length* of an admissible curve  $\xi$  is the integral  $\int_0^1 \left| \frac{d\xi}{dt} \right| dt$ , where  $|v| = \sqrt{\langle v | v \rangle_q} \quad \forall v \in T_q M$ . The infimum of the lengths of admissible curves connecting two points is the *Carnot–Caratheodory distance* between these points.

An important class of the sub-Riemannian structures is provided by magnetic fields on Riemannian surfaces. In this case  $M$  is the total space of a principal  $\mathbb{U}(1)$ -bundle over a Riemannian surface  $N$  and  $\Delta$  is a connection on the principal bundle. In other words,  $\Delta$  is a transversal to fibers  $\mathbb{U}(1)$ -invariant rank 2 distribution on  $M$ . The distribution is contact if and only if the curvature of the connection doesn't vanish. The inner product of a pair of vectors in  $\Delta_q$  equals the scalar product of their projections in  $T_q N$ .

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The length minimization problem for admissible curves is equivalent to the least action principle for the charged particles in the magnetic field (see [5] for details). The special case of the *constant* magnetic field (when the curvature of the connection is the area form multiplied by a constant) is locally equivalent to the classical Dido isoperimetric problem on the Riemannian surface.

The local structure of the Carnot–Caratheodory metric was studied in detail in the papers [1, 2, 3, 4]; it was shown that this structure is controlled by some fundamental differential invariants. In this note the same local invariants serve to express a global one: the Euler characteristic class of  $\Delta$ .

**2.** We assume that  $\Delta$  is an oriented contact structure. Then there exists a unique contact form  $\omega$  on  $M$  such that  $\Delta$  is the annihilator of  $\omega$ ,  $\Delta_q = \omega_q^\perp \forall q \in M$  and the form  $d\omega|_{\Delta_q}$  coincides with the area form on the oriented Euclidean plane  $\Delta_q$ . Let  $X_0$  be the characteristic (or Reeb) vector field of  $\omega$ ; it is defined by the relations  $X_0 \lrcorner d\omega = 0$ ,  $X_0 \lrcorner \omega = 1$ .

Let  $\delta(q_0, q_1)$  be the Carnot–Caratheodory distance between  $q_0$  and  $q_1$  and

$$\mathcal{C}_{q_0} = \{q_1 \in M : \text{the function } q \mapsto \delta(q_0, q) \text{ is not } C^1 \text{ at } q_1\}.$$

One can show that the line  $\mathbb{R}X_0(q)$  is the tangent cone<sup>1</sup> to the  $\mathcal{C}_q$  at  $q$ ,  $\mathbb{R}X_0(q) = T_q\mathcal{C}_q$ .

The curve  $t \mapsto e^{tX_0}(q)$  is transversal to the distribution  $\Delta$ . Let  $N_q \subset M$  be (a germ of) a 2-dimensional submanifold such that  $\Delta_q = T_qN$ ; then a neighborhood of  $q$  is sliced by the submanifolds  $e^{tX_0}(N_q)$ . One of the basic invariants of the sub-Riemannian structure arises from the asymptotics of the distance between  $q$  and  $e^{tX_0}(N_q) \cap \mathcal{C}_q$  as  $t \rightarrow 0$ :

$$\delta(q, e^{tX_0}(N_q) \cap \mathcal{C}_q) = (2 - \kappa(q)t)|\pi t|^{1/2} + O(t^2),$$

where  $\kappa$  is a smooth function on  $M$ .

Let  $X_1(q), X_2(q)$  be an orthonormal frame in  $\Delta_q$ . The explicit expression of  $\kappa$  in terms of the moving frame  $X_0, X_1, X_2$  is as follows:

$$\kappa = X_1 c_{12}^2 - X_2 c_{12}^1 - (c_{12}^1)^2 - (c_{12}^2)^2 + \frac{1}{2}(c_{02}^1 - c_{01}^2),$$

where  $[X_i, X_j] = \sum_k c_{ij}^k X_k$ .

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<sup>1</sup>There are many definitions of the tangent cone to the closed set; all of them provide one and the same cone in this particular case.

In the special case of the constant magnetic field on the Riemannian surface (see Sec. 1)  $X_0$  is a generator of the structural group of the principle bundle,  $\kappa$  is constant on the fibers of the bundle and is actually the pullback of the Gaussian curvature of the Riemannian surface.

**3.** Now we take the dual object to the sub-Riemannian structure, the Hamiltonian  $h$  on the cotangent bundle  $T^*M$ :

$$h(\lambda) = \frac{1}{2}(\max\{\langle \lambda, v \rangle : v \in \Delta_q, |v| = 1\})^2, \quad \lambda \in T_q^*M, q \in M.$$

This Hamiltonian serves to describe sub-Riemannian geodesics, i.e. admissible curves whose small pieces are length minimizers. It follows from the Pontryagin Maximum Principle that geodesics are exactly projections to  $M$  of the trajectories of the Hamiltonian system in  $T^*M$  associated with  $h$ .

Let  $u_i(\lambda) = \langle \lambda, X_i(q) \rangle$ ,  $i = 1, 2$ ; then  $h(\lambda) = \frac{1}{2}(u_1(\lambda)^2 + u_2(\lambda)^2)$ ,  $\lambda \in T_q^*M$ . The Hamiltonian keeps all the information on the sub-Riemannian structure: both  $\Delta$  and the inner product are easily recovered from  $h$ .

We denote:  $u_0(\lambda) = \langle \lambda, X_0(q) \rangle$ ,  $\lambda \in T_q^*M$ ,  $q \in M$ , and  $\Delta^* = u^{-1}(0)$ . Then  $\Delta^*$  is a rank 2 linear subbundle of  $T^*M$  with the fibers  $\Delta_q^* = \Delta^* \cap T_q^*M$ . Obviously,  $(\lambda, \xi) \mapsto \langle \lambda, \xi \rangle$ ,  $\lambda \in \Delta_q^*$ ,  $\xi \in \Delta_q$ , is a nondegenerate pairing. Moreover, quadratic form  $2h|_{\Delta_q^*}$  defines the Euclidean structure on  $\Delta_q^*$  dual to the given Euclidean structure on  $\Delta_q$ .

We'll deal with homogeneous polynomials on the plane  $\Delta_q^*$ . The group  $SO(\Delta_q^*)$  acts on the polynomials by the changing of variables. Irreducible components of this (real) action are 2-dimensional spaces of polynomials having the following expression in polar coordinates  $(r, \theta)$ :

$$\text{span}\{r^n \cos(k\theta), r^n \sin(k\theta)\}, \quad (1)$$

where  $n - k$  is a nonnegative even number. Note that  $r^2 = 2h|_{\Delta_q^*}$ . So any homogeneous degree  $n$  polynomial  $\phi$  has a unique presentation as a sum of isotopic components:  $\phi = \sum_{i=0}^{[n/2]} \phi^{n-2i}$ , where  $\phi^k$  belongs to the space (1). This presentation is actually equivalent to the Fourier expansion of the restriction of  $\phi$  to the unit circle.

**4.** Recall that  $h, u_0$  are functions on the cotangent bundle  $T^*M$ , where restrictions of  $h$  to the fibers  $T_q^*M$  are quadratic forms and restrictions of  $u_0$  to the fibers are linear forms. Hence the Poisson bracket  $\{h, u_0\}$  is quadratic and the double Poisson bracket  $\{h, \{h, u_0\}\}$  is cubic on the fibers. Let  $\phi_q = \{h, \{h, u_0\}\}|_{\Delta_q^*}$ ; then  $\phi_q = \phi_q^1 + \phi_q^3$ , where  $\phi_q^1$  is the product of  $h|_{\Delta_q^*}$  and a

linear form on  $\Delta_q^*$ , according to (1). In other words,  $\phi_q^1(\lambda) = h(\lambda)\langle\lambda, f(q)\rangle$  for some  $f(q) \in \Delta_q$ .

We thus obtain an intrinsically defined vector field  $f$  with values in  $\Delta$  in addition to the transversal to  $\Delta$  field  $X_0$ .

**Theorem 1** *The 2-form on  $M$*

$$\left( \frac{\kappa}{2\pi} X_0 - \frac{1}{\pi} f \right) \Big| \omega \wedge d\omega \quad (2)$$

*is closed and represents the Euler characteristic class of the oriented linear bundle  $\Delta$ .*

**Remark 1.** The statement of the theorem can be also formulated as follows: The flow generated by the field

$$\left( \frac{\kappa}{2\pi} X_0 - \frac{1}{\pi} f \right) \quad (3)$$

preserves the volume form  $\omega \wedge d\omega$  and the *rotation class* (see [6]) of the field (3) is equal to the Euler class of  $\Delta$ .

**Remark 2.** In the case of the principal bundle  $M \xrightarrow{\text{U}(1)} N$  and the sub-Riemannian structure defined by a *constant* magnetic field (see sec.1) we have  $\{h, u_0\} = 0$  and hence  $f = 0$ . Then (2) takes the form :  $\frac{\kappa}{2\pi} X_0 \Big| \omega \wedge d\omega = \frac{\kappa}{2\pi} d\omega$ . Moreover,  $d\omega$  and  $\kappa$  are the pullbacks of the area form and of the Gauss curvature on  $N$  so that the form (2) turns into the pullback of the Gauss–Bonnet form on  $N$ .

The proof of Theorem 1 consists of a calculation with moving frames in  $T^*M$ . The idea is to construct an appropriate linear connection on the bundle  $\Delta^* \subset T^*M$  via the Hamiltonian vector fields associated with  $h$  and  $u_0$  and a vertical vector field generating rotations of the fibers  $\Delta_q^*$ . The form (2) is the curvature form of the correspondent linear connection.

## References

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