

Extragalactic γ -ray signal from dark matter annihilation: a power spectrum perspective

DSU2013, SISSA - October 16, 2013

Emiliano SEFUSATTI

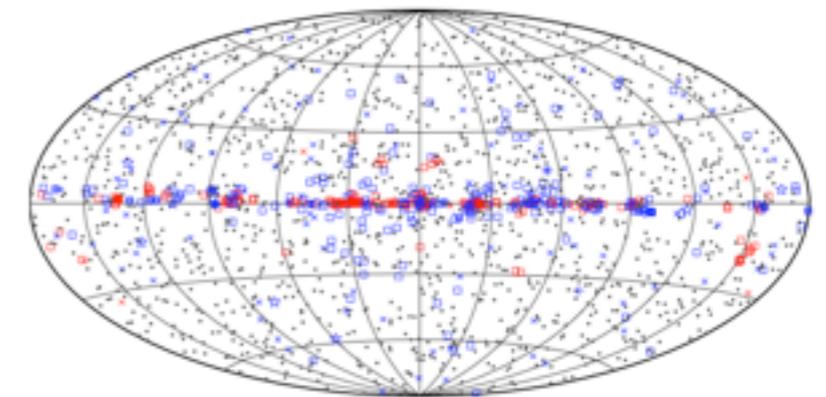
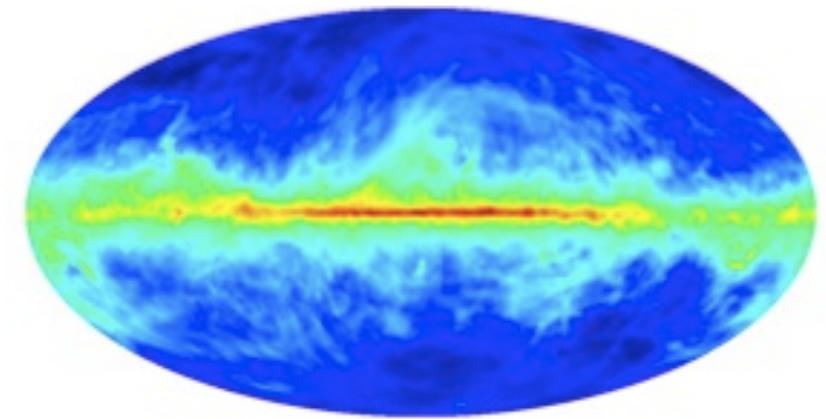
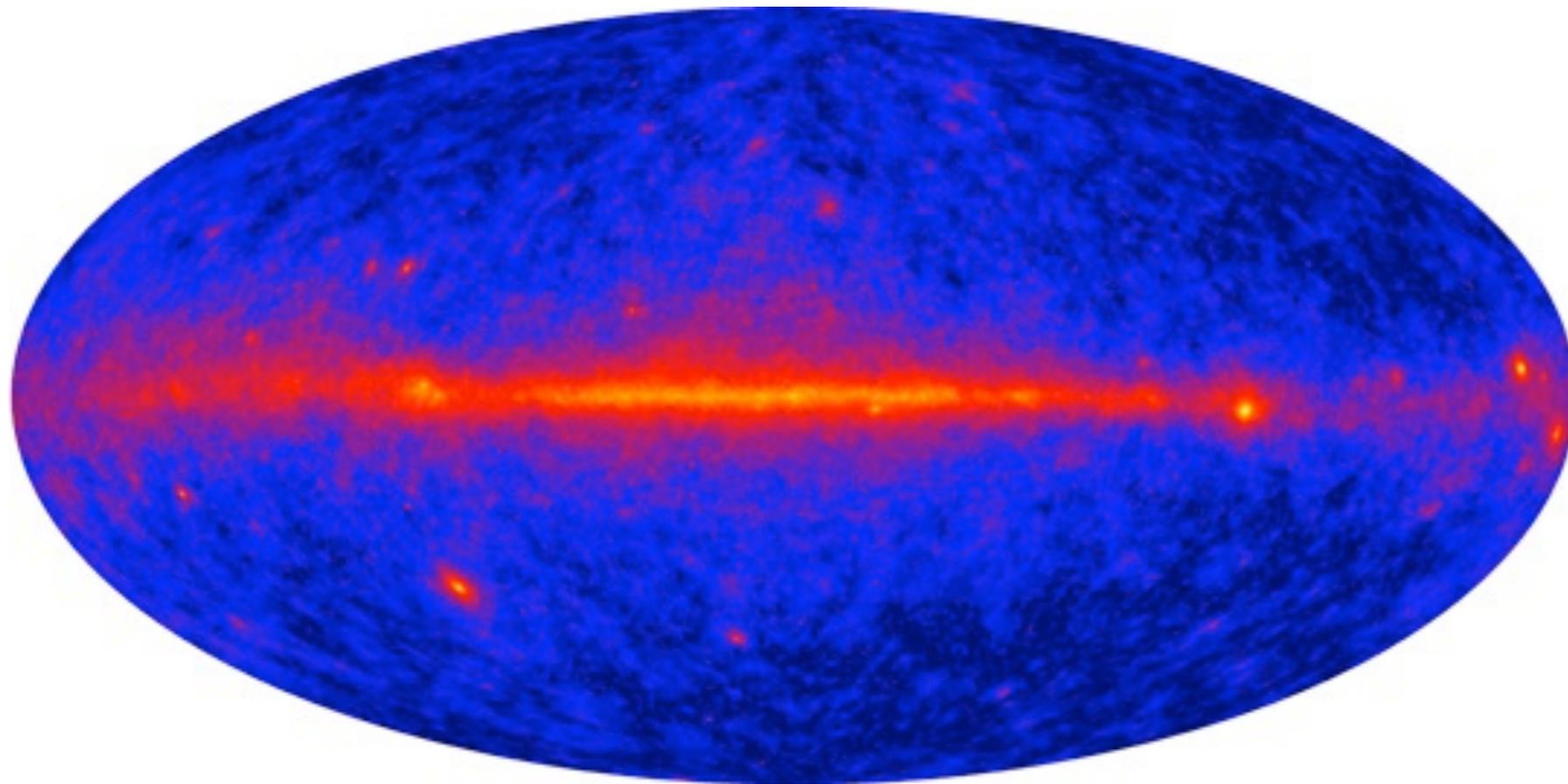
w/ G. Zaharijas, P. Serpico, M. Gustafsson

MNRAS **421**, L87 (2012) and *in preparation*

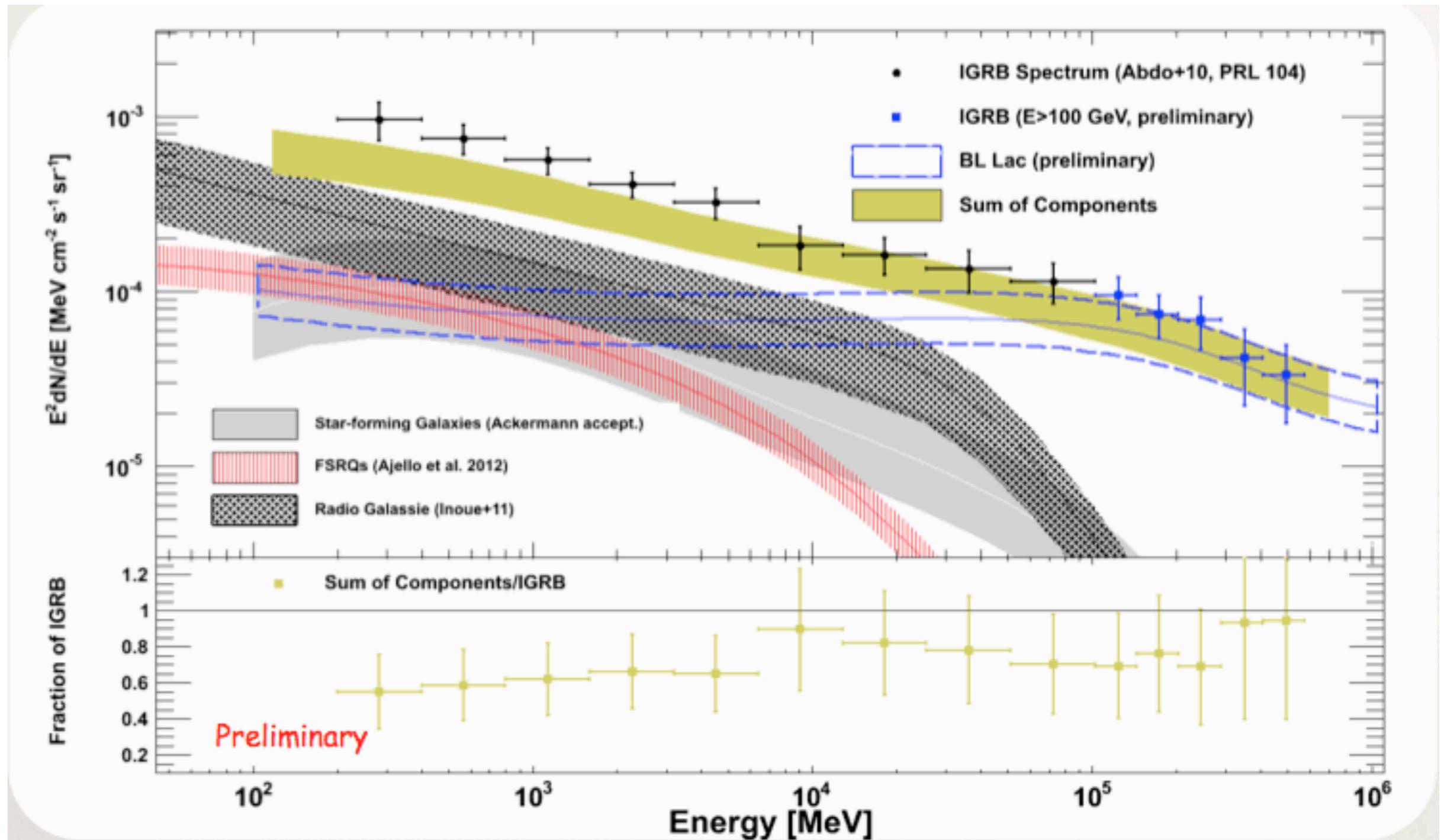


The γ -ray sky

The **Extra Galactic, diffuse γ -ray signal** is inferred from a multicomponent fit to Fermi-LAT data.



Extragalactic diffuse emission: DM contribution?



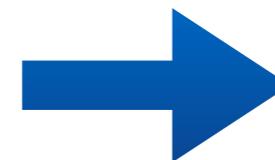
Sources:

- **unresolved extragalactic sources**: blazars, star forming and star burst galaxies ...
- **Dark Matter annihilation** in all halos at all red-shifts can contribute

Extragalactic γ -rays from annihilating DM

$$n = \frac{\langle \sigma v \rangle}{2 m_{DM}^2} \int_V d^3 x \rho_{DM}^2(\vec{x})$$

of annihilations per unit volume



$$\frac{d\phi_\gamma}{dE_0} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{m_{DM}^2} \int dz \frac{(1+z)^3}{H(z)} \frac{dN_\gamma[E_0(1+z)]}{dE} e^{-\tau(z, E_0)} \rho^2(z)$$

observed spectrum



DM “interesting properties”



geometry



spectrum at emission



attenuation due to pair production on extragalactic background light



matter density

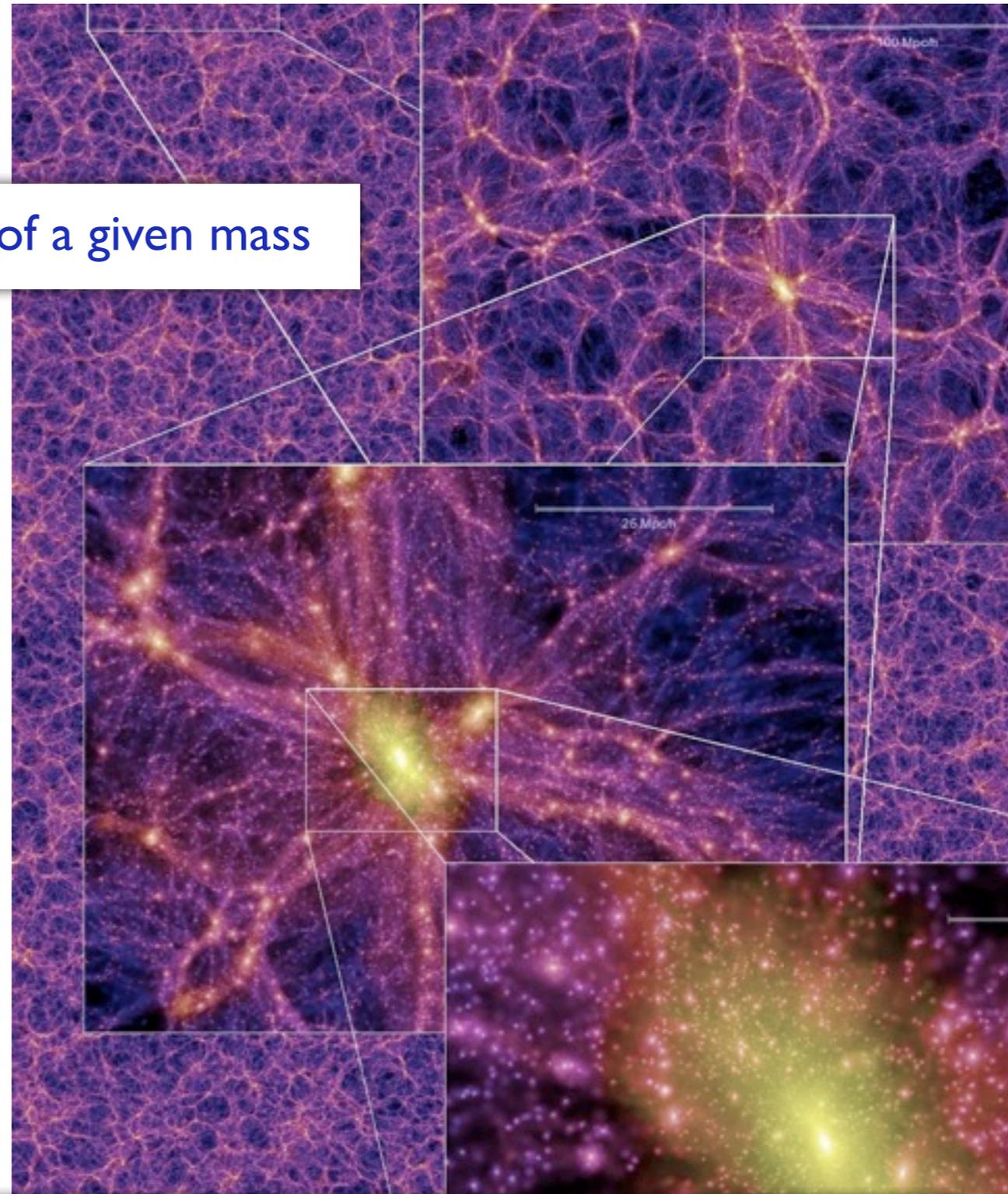
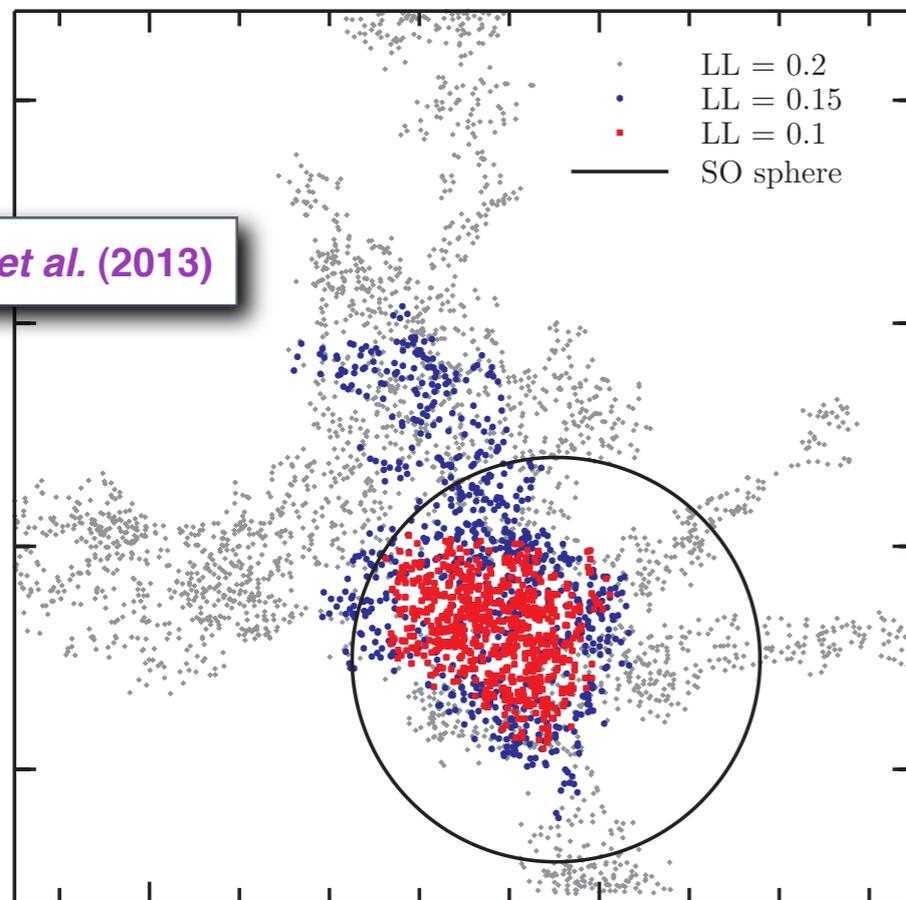


The Halo Model approach

$$\rho(z, \hat{\Omega}) = \bar{\rho}(z) \Delta(z, \hat{\Omega})$$

all DM particles belong to a (spherical?) DM *halo* of a given mass

Watson et al. (2013)



then have the “flux multiplier”

$$\Delta^2(z) \equiv \zeta(z) = \frac{1}{\Omega_m \rho_c} \int_{M_{min}} dM \frac{dn}{dM} M \frac{\Delta_{vir}(z)}{3} \langle F \rangle \quad F = c_v^3(M, z) \frac{\int_0^{c_v} dx x^2 \kappa^2(x)}{\left[\int_0^{c_v} dx x^2 \kappa(x) \right]^2}$$

The Halo Model approach

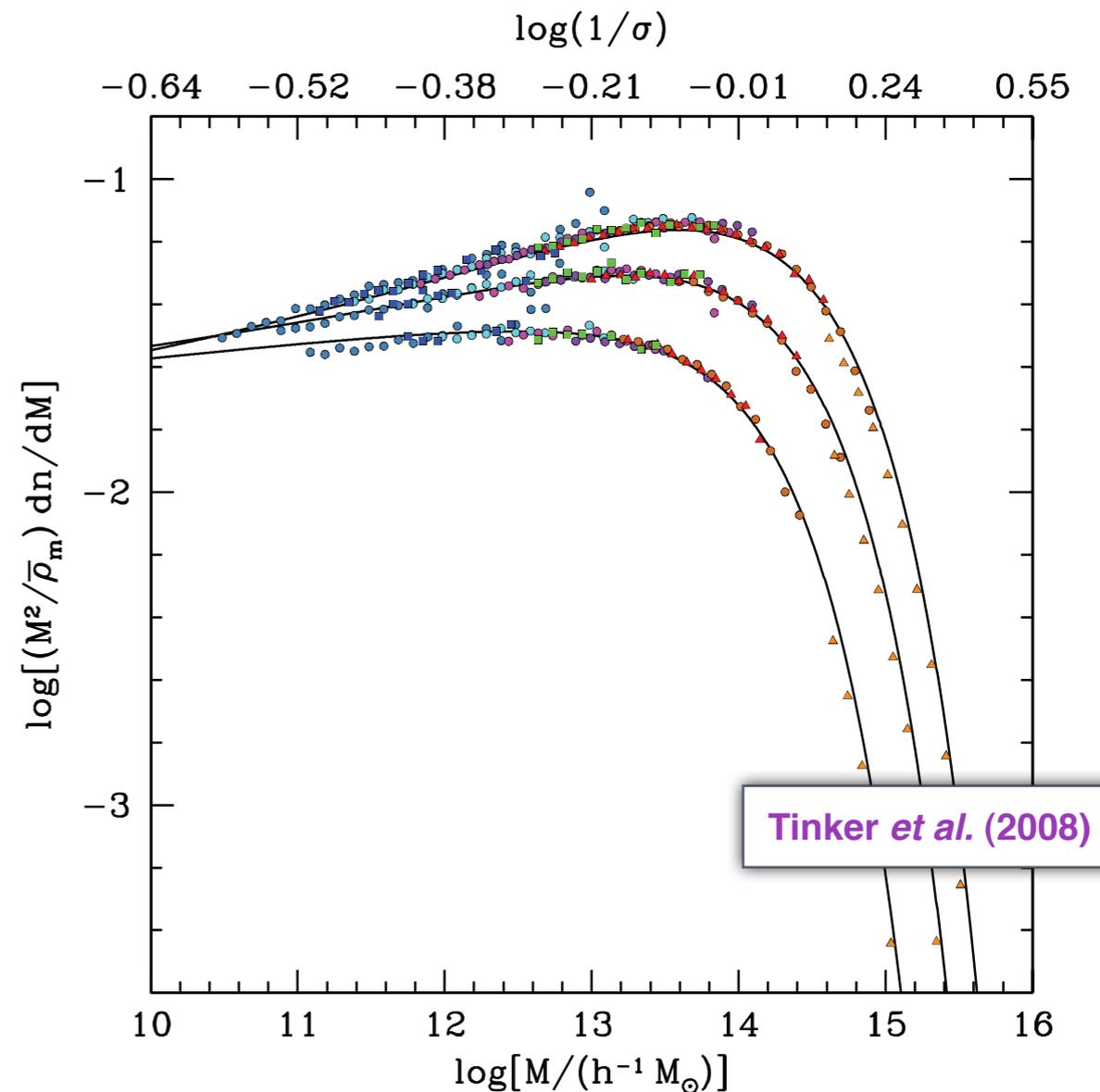
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Ingredients:

1. Halos **mass function**



The Halo Model approach

$$\rho(z, \hat{\Omega}) = \bar{\rho}(z) \Delta(z, \hat{\Omega})$$

all DM particles belong to a DM *halo* of a given mass

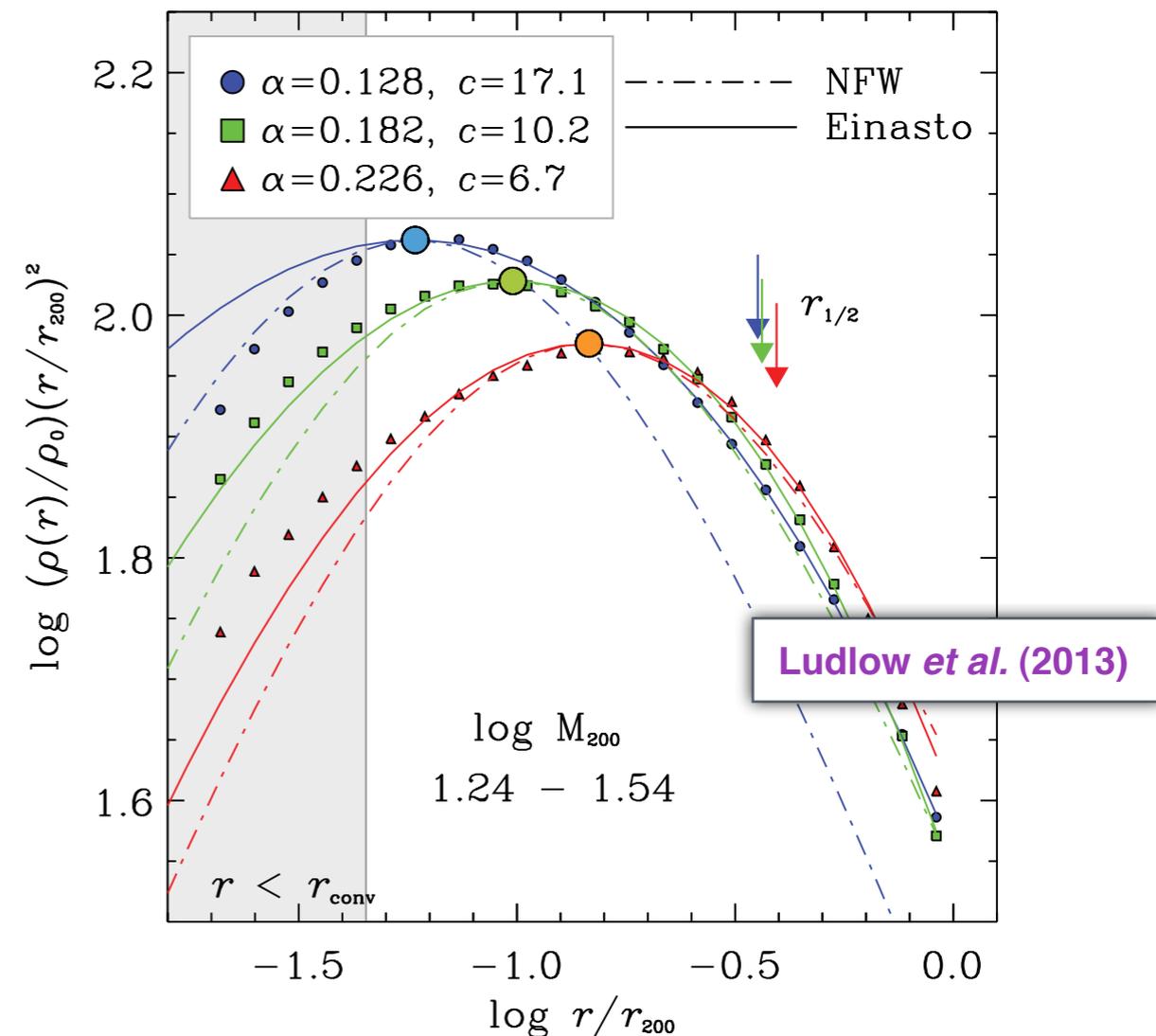
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Ingredients:

1. Halos mass function
2. Halos **density profile** (NFW, Einasto, etc ...)

e.g.

$$\rho_{NFW}(r) = \frac{\rho_0}{\frac{r}{R_s} \left(1 + \frac{r}{R_s} \right)^2}$$



The Halo Model approach

$$\rho(z, \hat{\Omega}) = \bar{\rho}(z) \Delta(z, \hat{\Omega})$$

all DM particles belong to a DM *halo* of a given mass

$$\Delta^2(z) \equiv \zeta(z) = \frac{1}{\Omega_m \rho_c} \int_{M_{min}} dM \frac{dn}{dM} M \frac{\Delta_{vir}(z)}{3} \langle F \rangle$$

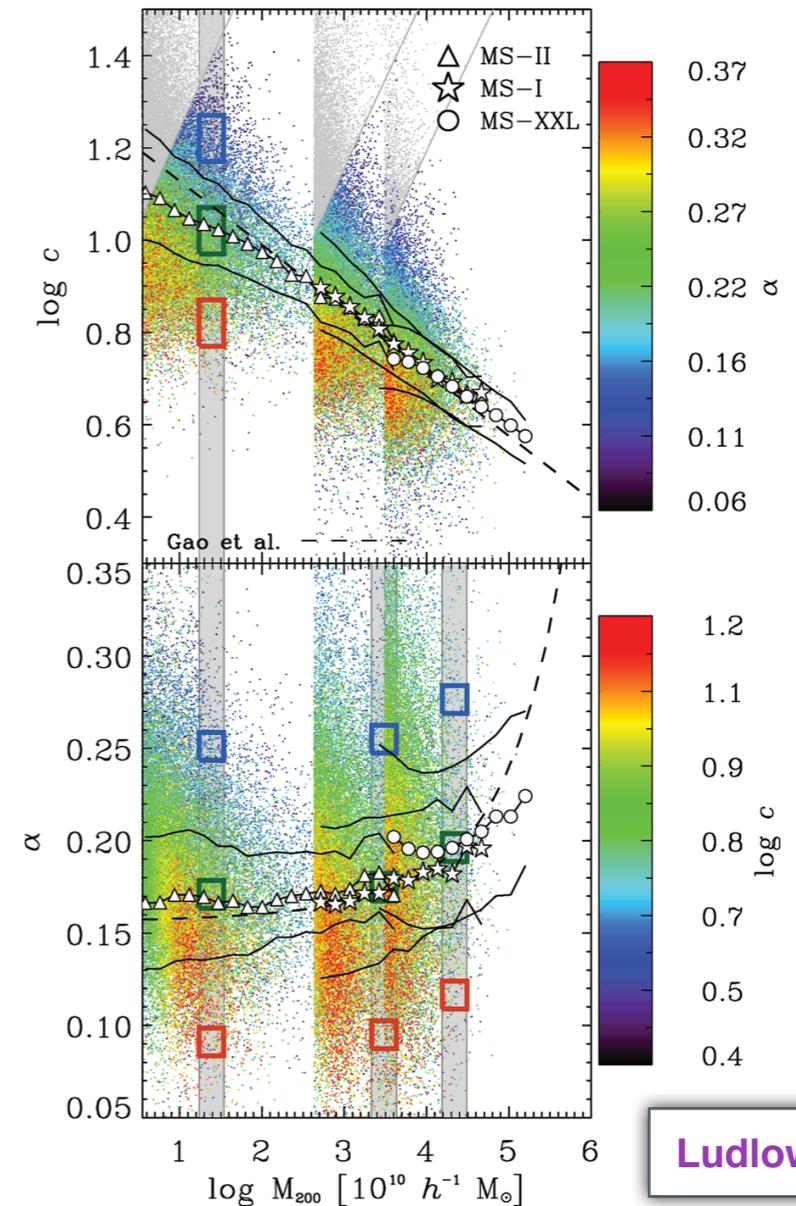
$$F = c_v^3(M, z) \frac{\int_0^{c_v} dx x^2 \kappa^2(x)}{[\int_0^{c_v} dx x^2 \kappa(x)]^2}$$

Ingredients:

1. Halos mass function
2. Halos density profile (NFW, Einasto, etc ...)
3. Halos **concentration**

e.g.

$$c_v(M, z) = \frac{R_v}{R_s}$$



Ludlow et al. (2013)

The Halo Model approach

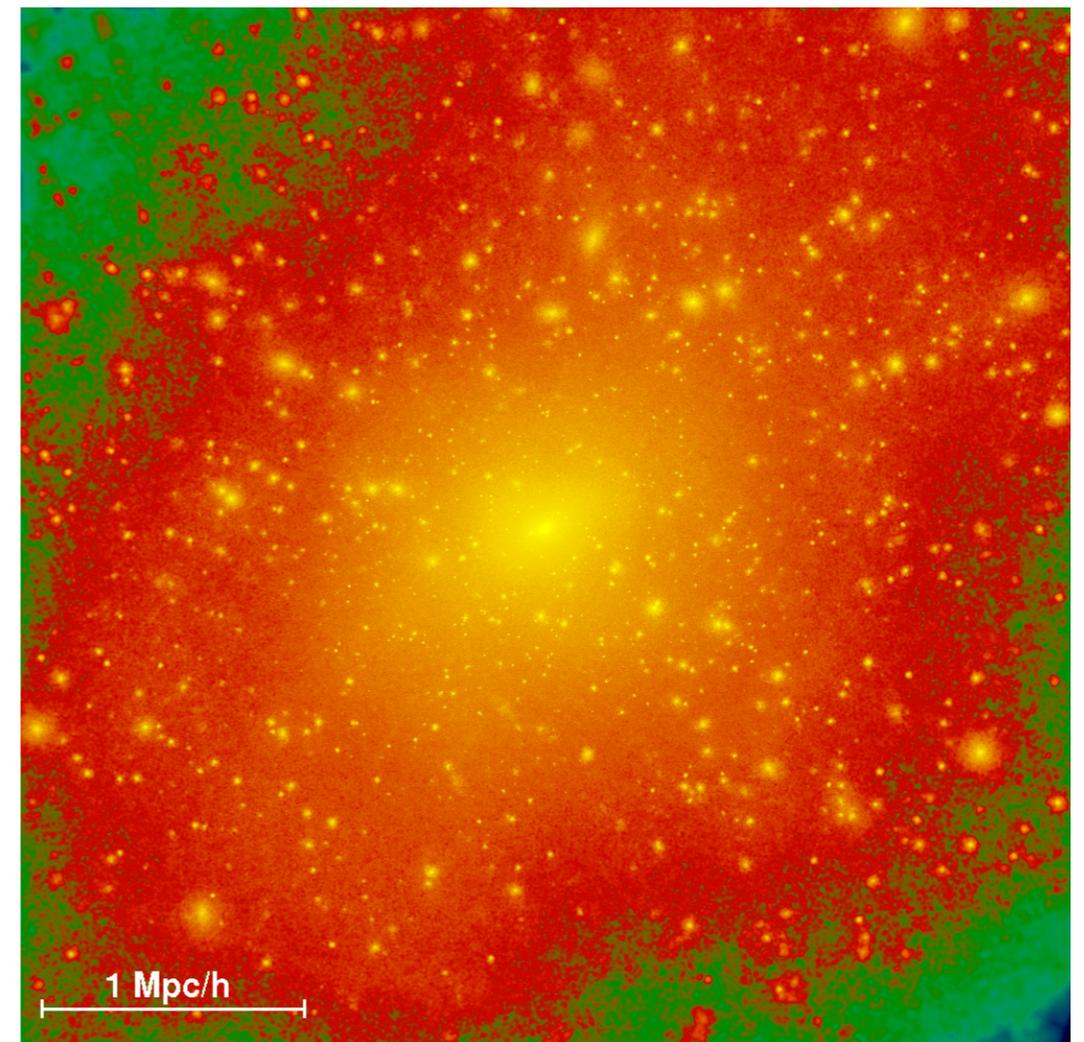
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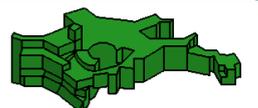
Ingredients:

1. Halos mass function
 2. Halos density profile (NFW, Einasto, etc ...)
 3. Halos concentration
- + all of the above for **subhalos**



Dark matter in the S3 cluster.

Springel, White & Tormen (1999)



The Halo Model approach

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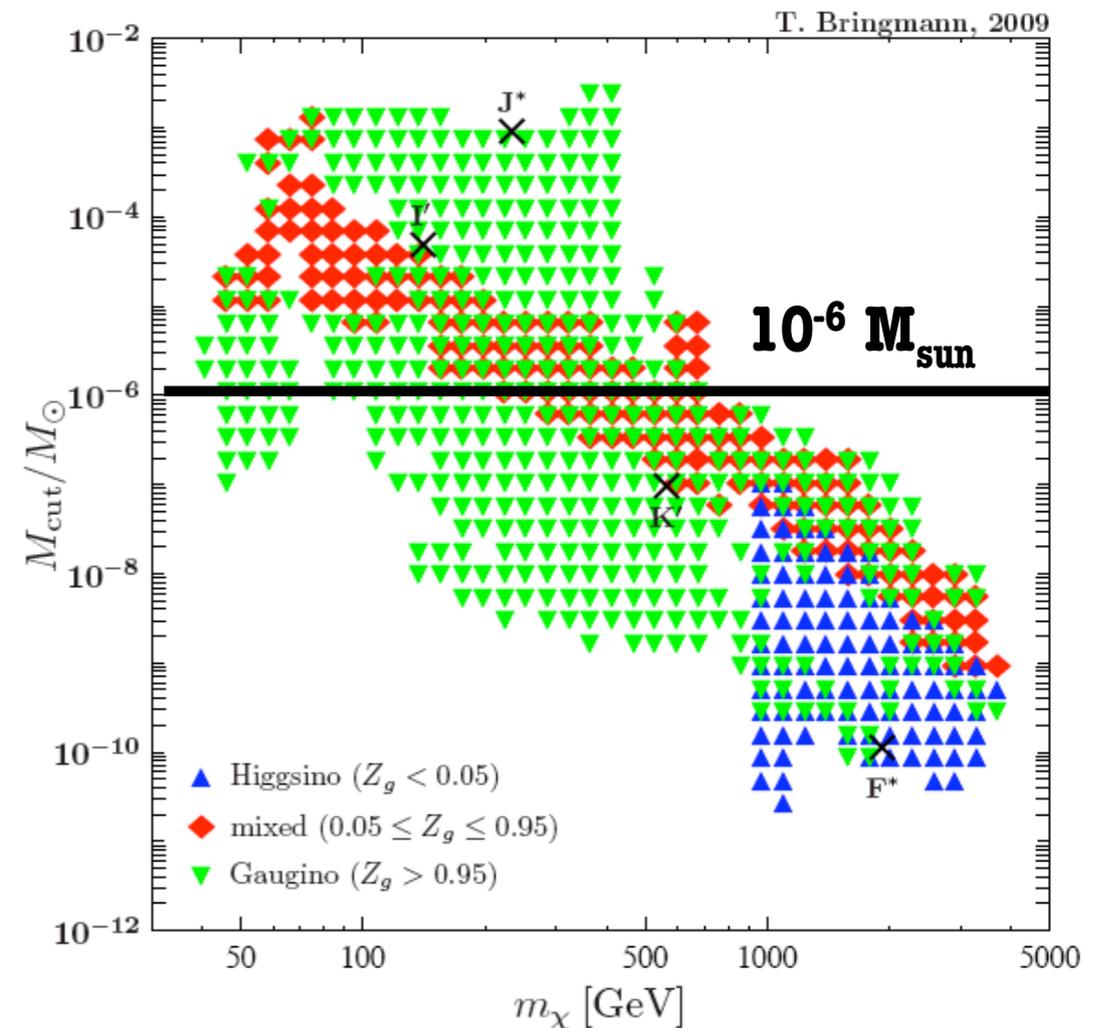
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Typical M_{min} for a WIMP = $10^{-6} M_{sun}$



The Halo Model approach

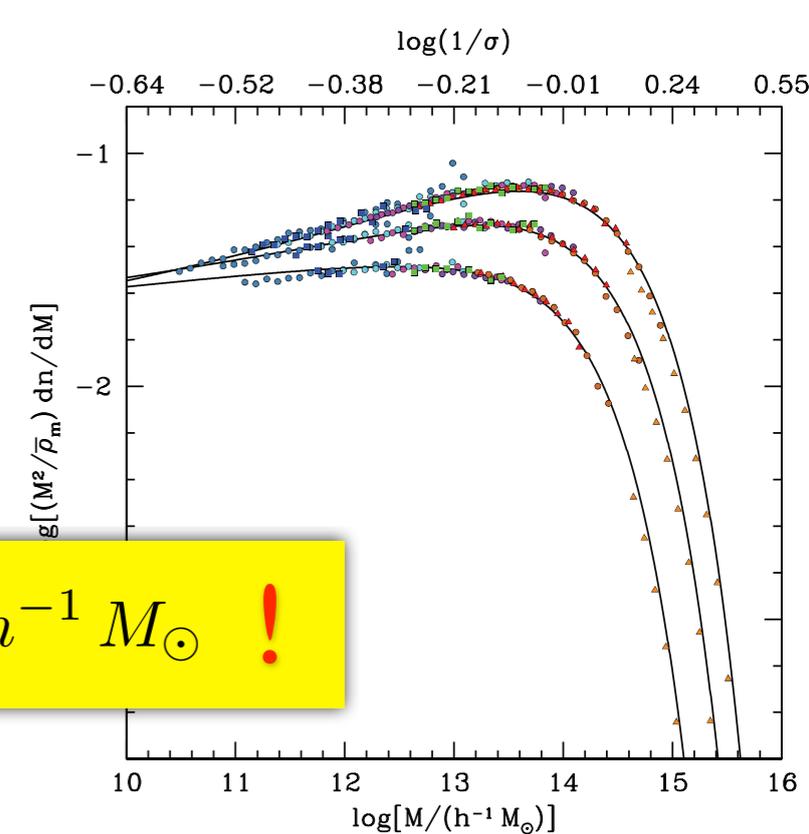
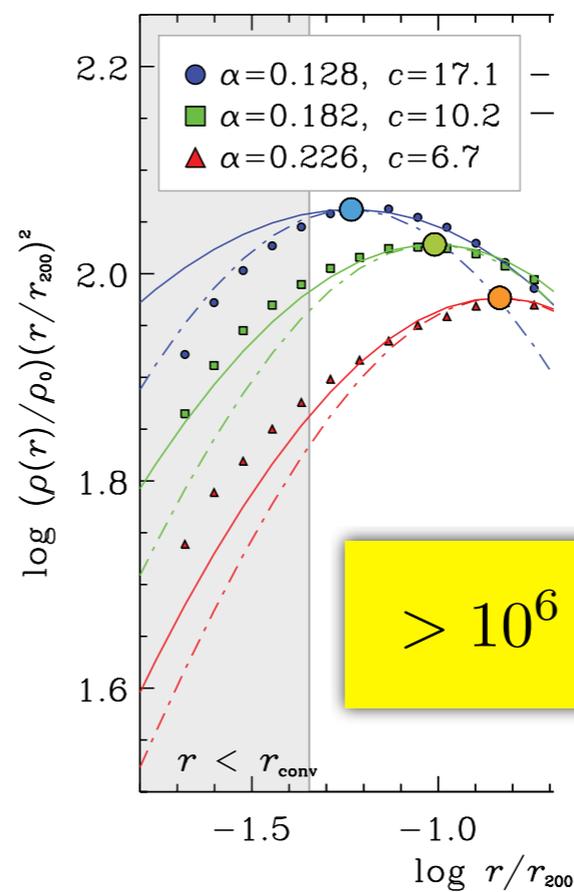
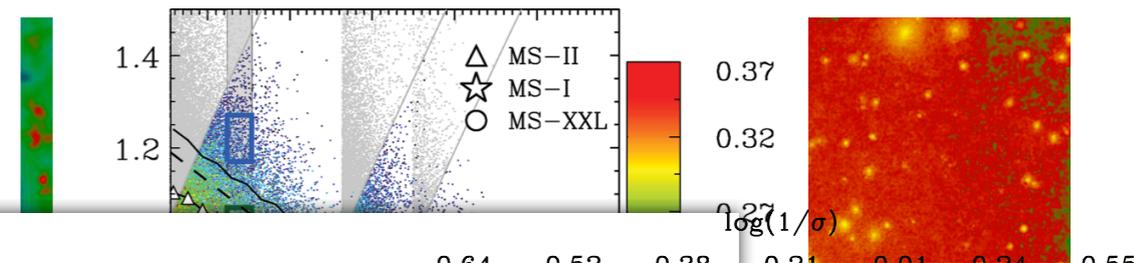
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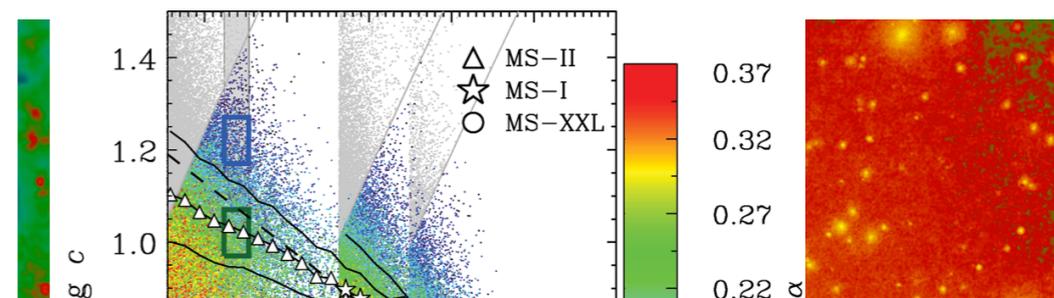
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an extrapolation of more than 10 orders of magnitude!

$> 10^6 h^{-1} M_{\odot}$

$10^{-9}, 10^{-6} h^{-1} M_{\odot} ?$



11 12 13 14 15 16
log[M/(h⁻¹ M_⊙)]

A Power Spectrum approach

$$\rho(z, \hat{\Omega}) = \bar{\rho}(z) \Delta(z, \hat{\Omega}) = \bar{\rho}(z) [1 + \delta(z, \hat{\Omega})] \longrightarrow \zeta(z) = \langle \delta^2(z, \hat{\Omega}) \rangle$$

Enter: the dark matter power spectrum

$$\langle \delta(\vec{x}) \delta(\vec{y}) \rangle \equiv \xi(|\vec{x} - \vec{y}|) = \frac{1}{(2\pi)^3} \int e^{ik|\vec{x} - \vec{y}|} P(k)$$

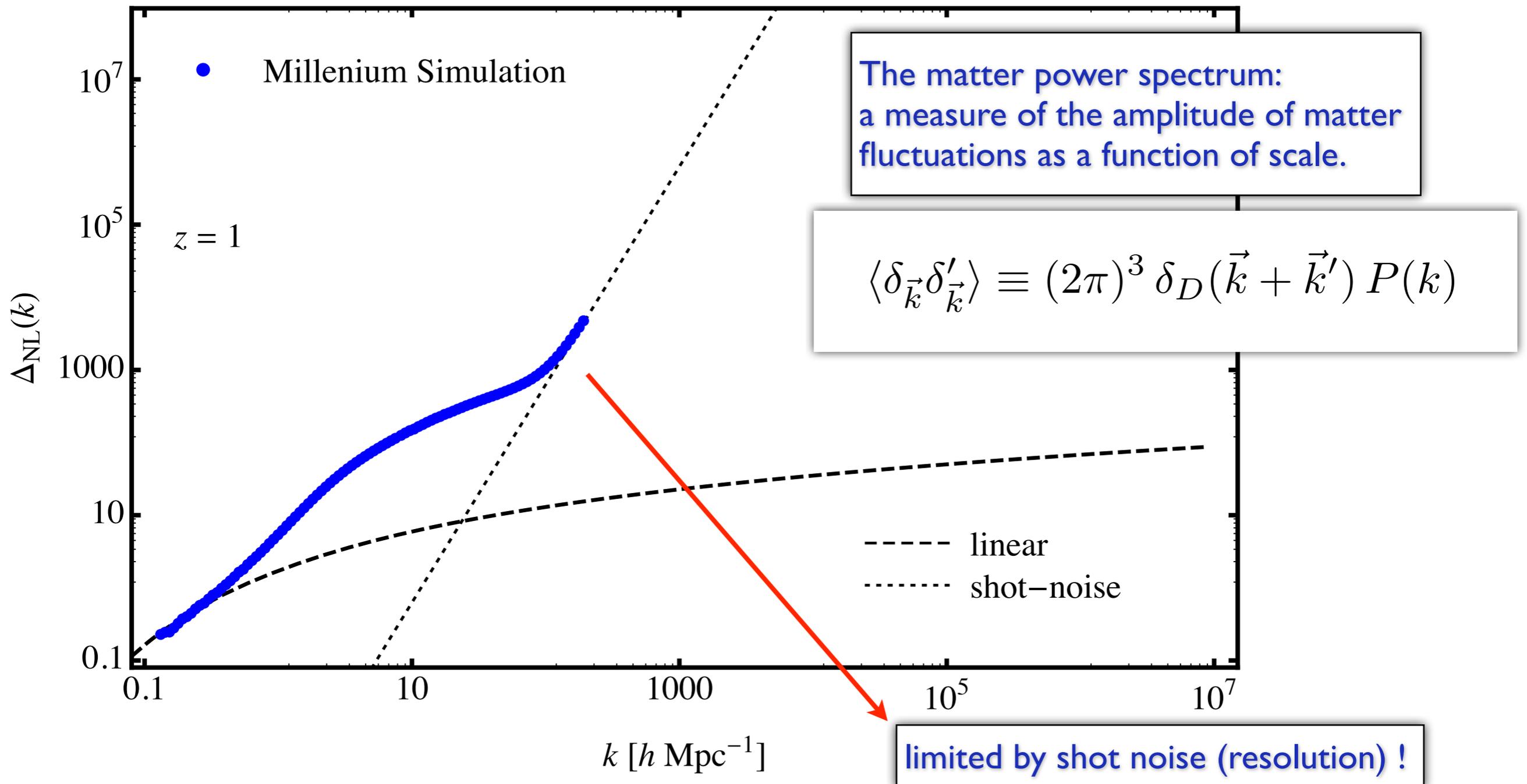
$$\langle \delta^2(\vec{x}) \rangle \equiv \xi(0) = \frac{1}{2\pi^2} \int dk k^2 P(k)$$

$$\zeta(z) = \langle \delta^2(z, \hat{\Omega}) \rangle = \int_0^{k_{max}} \frac{dk}{k} \frac{k^3 P_{NL}(k)}{2\pi^2}$$

all-in-one function!

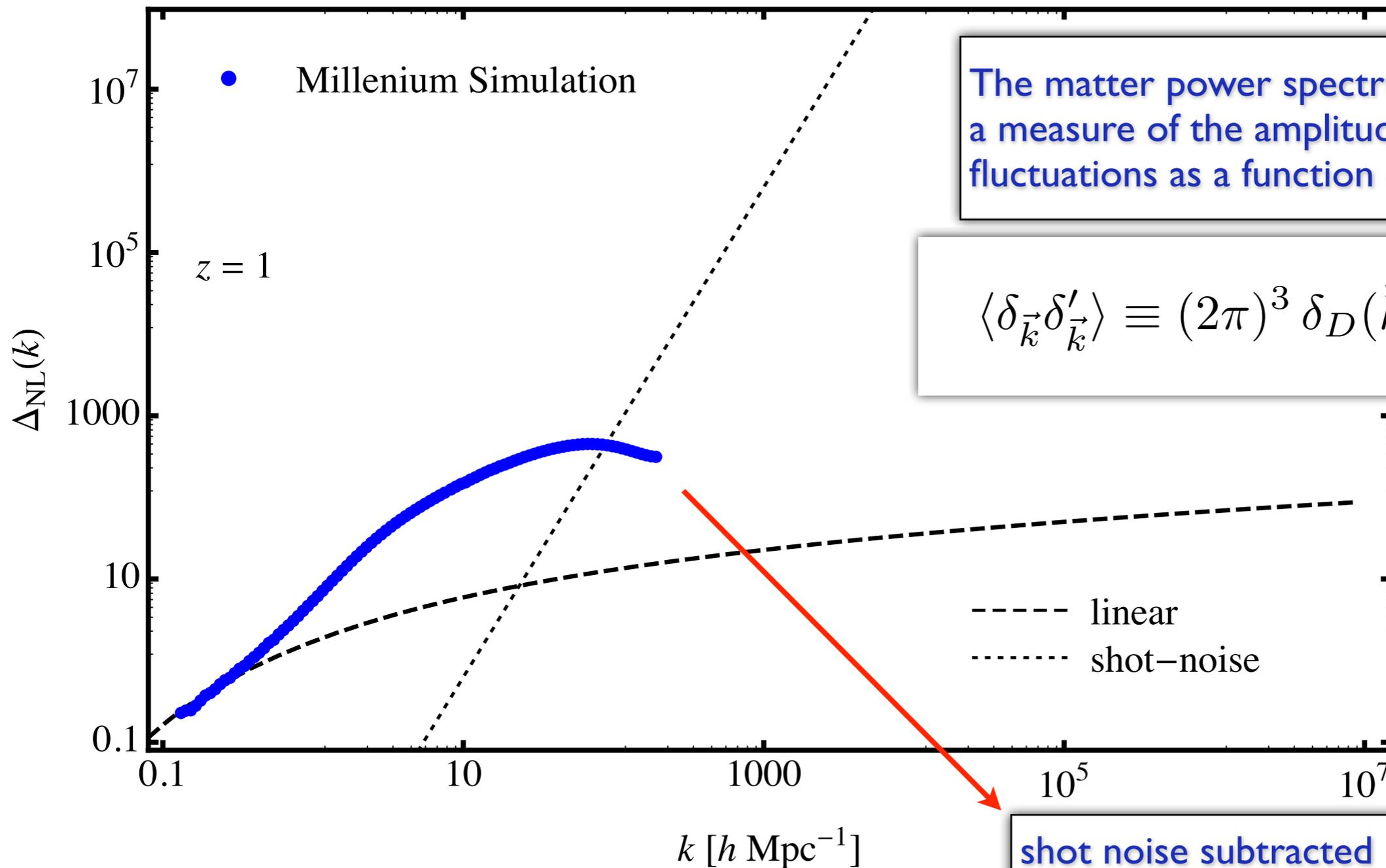
The nonlinear matter Power Spectrum

$$\zeta(z) = \langle \delta^2(z, \hat{\Omega}) \rangle = \int_0^{k_{max}} \frac{dk}{k} \Delta_{NL}(k) \quad \Delta_{NL}(k) \equiv \frac{k^3 P_{NL}(k)}{2\pi^2}$$



The nonlinear matter Power Spectrum

$$\zeta(z) = \langle \delta^2(z, \hat{\Omega}) \rangle = \int_0^{k_{max}} \frac{dk}{k} \Delta_{NL}(k) \quad \Delta_{NL}(k) \equiv \frac{k^3 P_{NL}(k)}{2\pi^2}$$



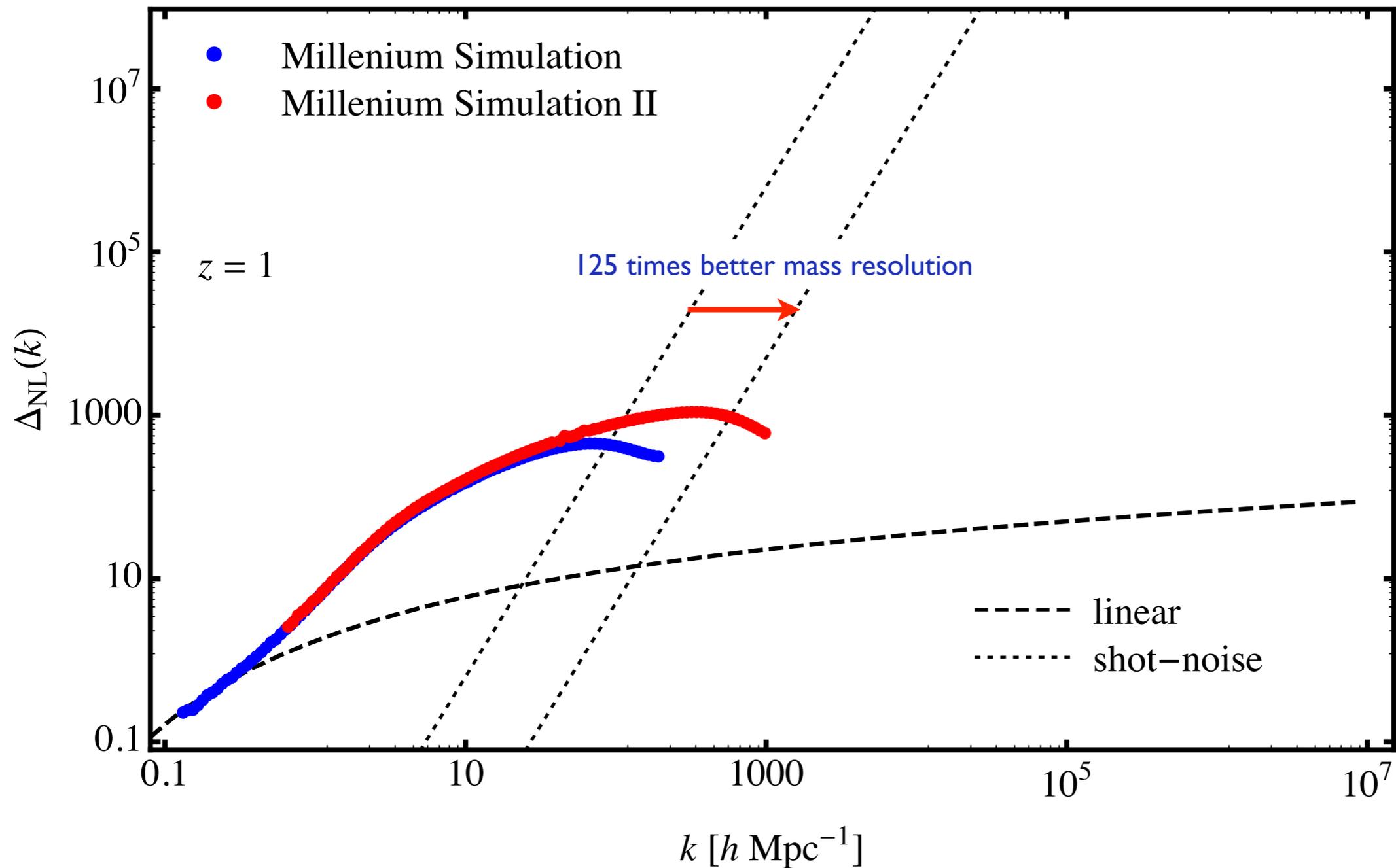
The matter power spectrum:
 a measure of the amplitude of matter
 fluctuations as a function of scale.

$$\langle \delta_{\vec{k}} \delta'_{\vec{k}'} \rangle \equiv (2\pi)^3 \delta_D(\vec{k} + \vec{k}') P(k)$$

shot noise subtracted

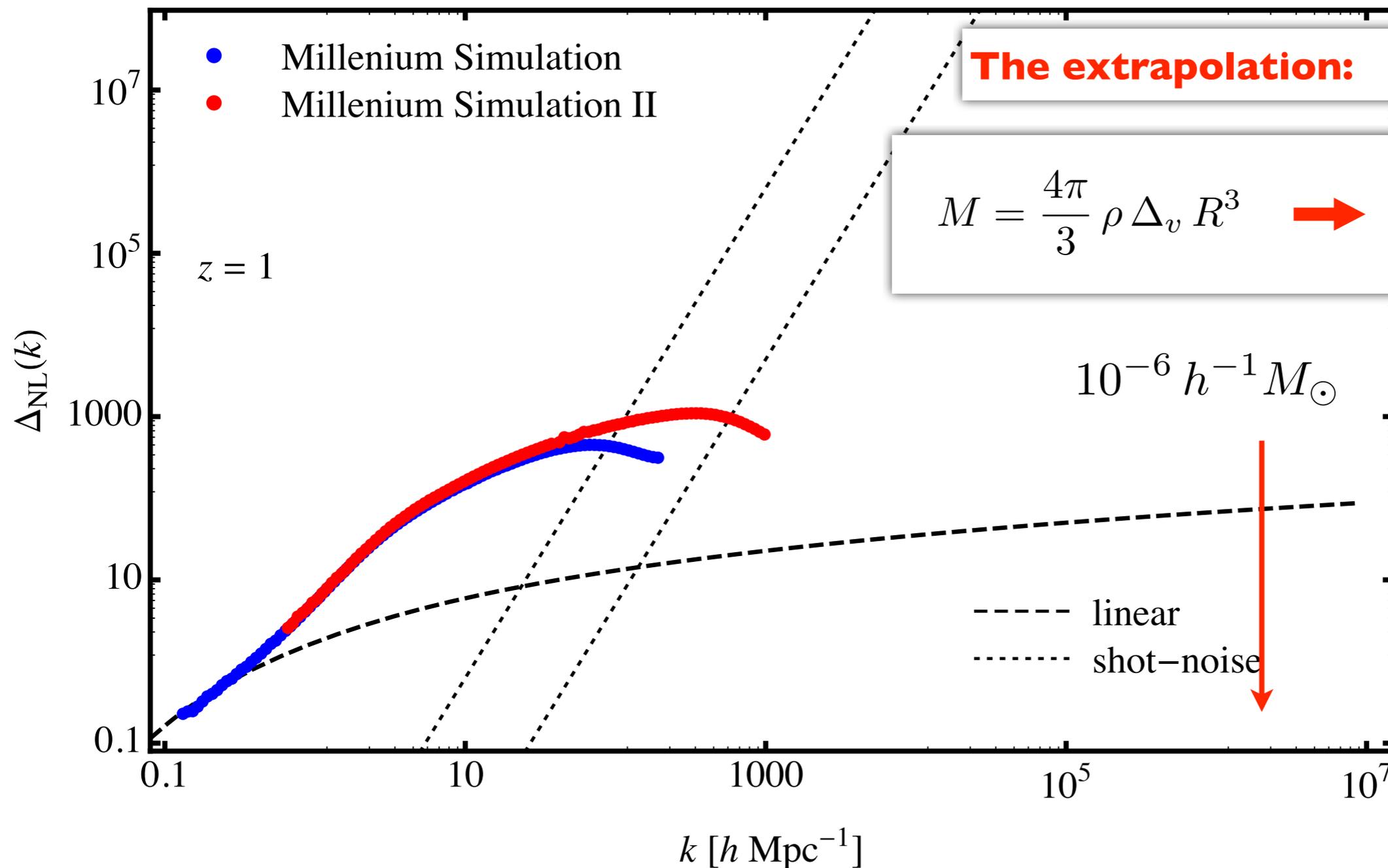
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The nonlinear matter Power Spectrum

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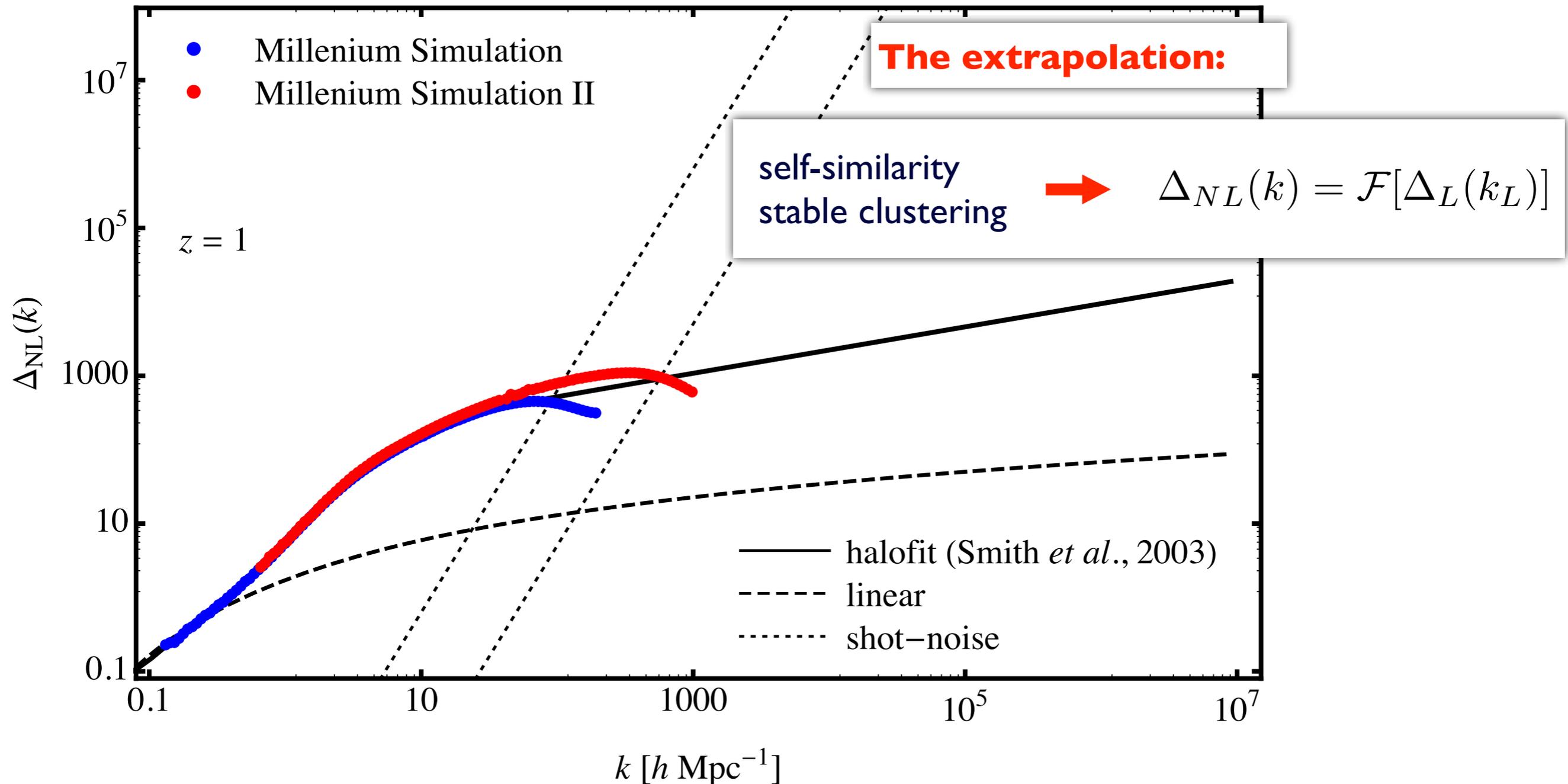


The extrapolation:

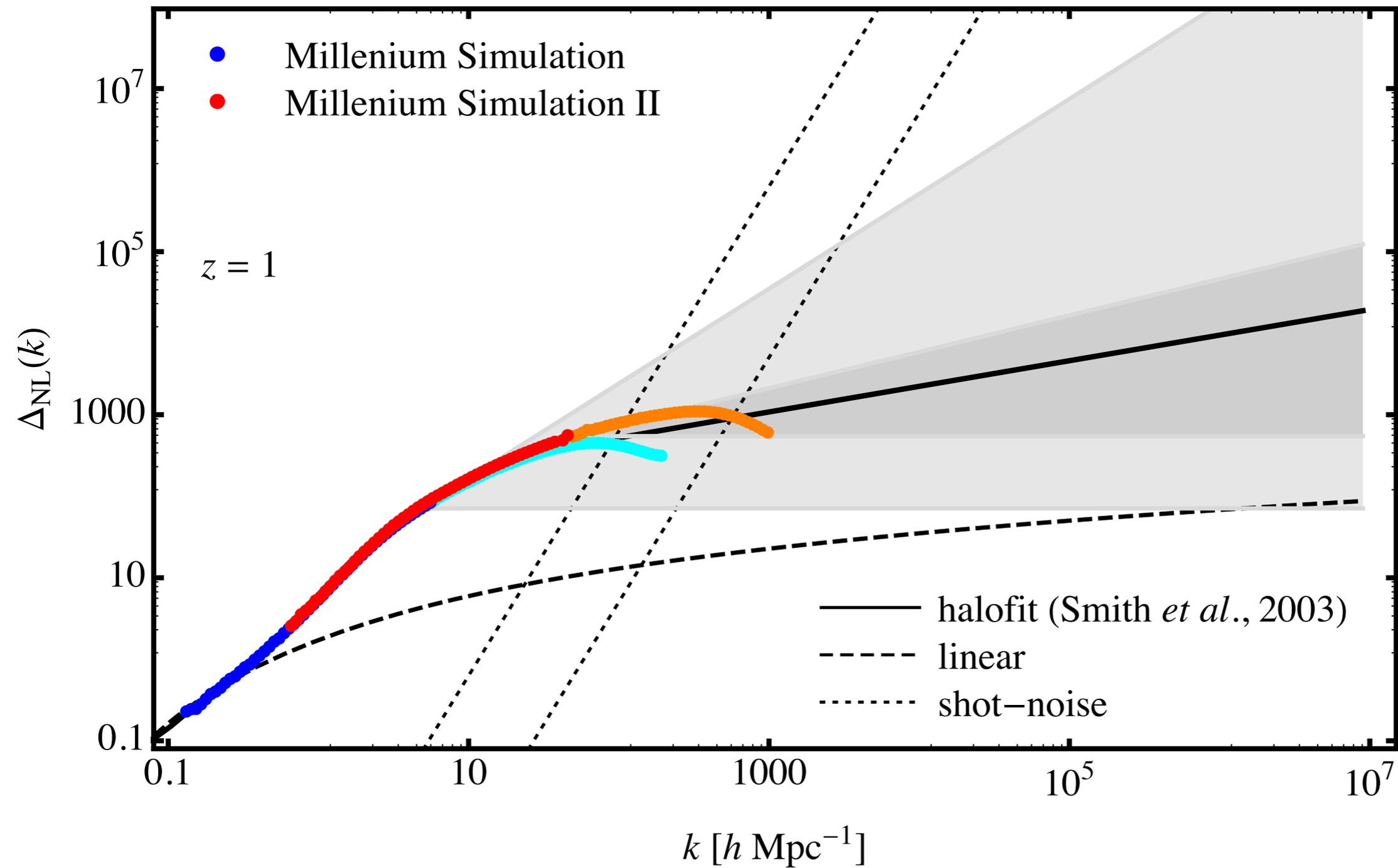
$$M = \frac{4\pi}{3} \rho \Delta_v R^3 \quad \rightarrow \quad k_{max} \sim M_{min}^{-1/3}$$

The nonlinear matter Power Spectrum

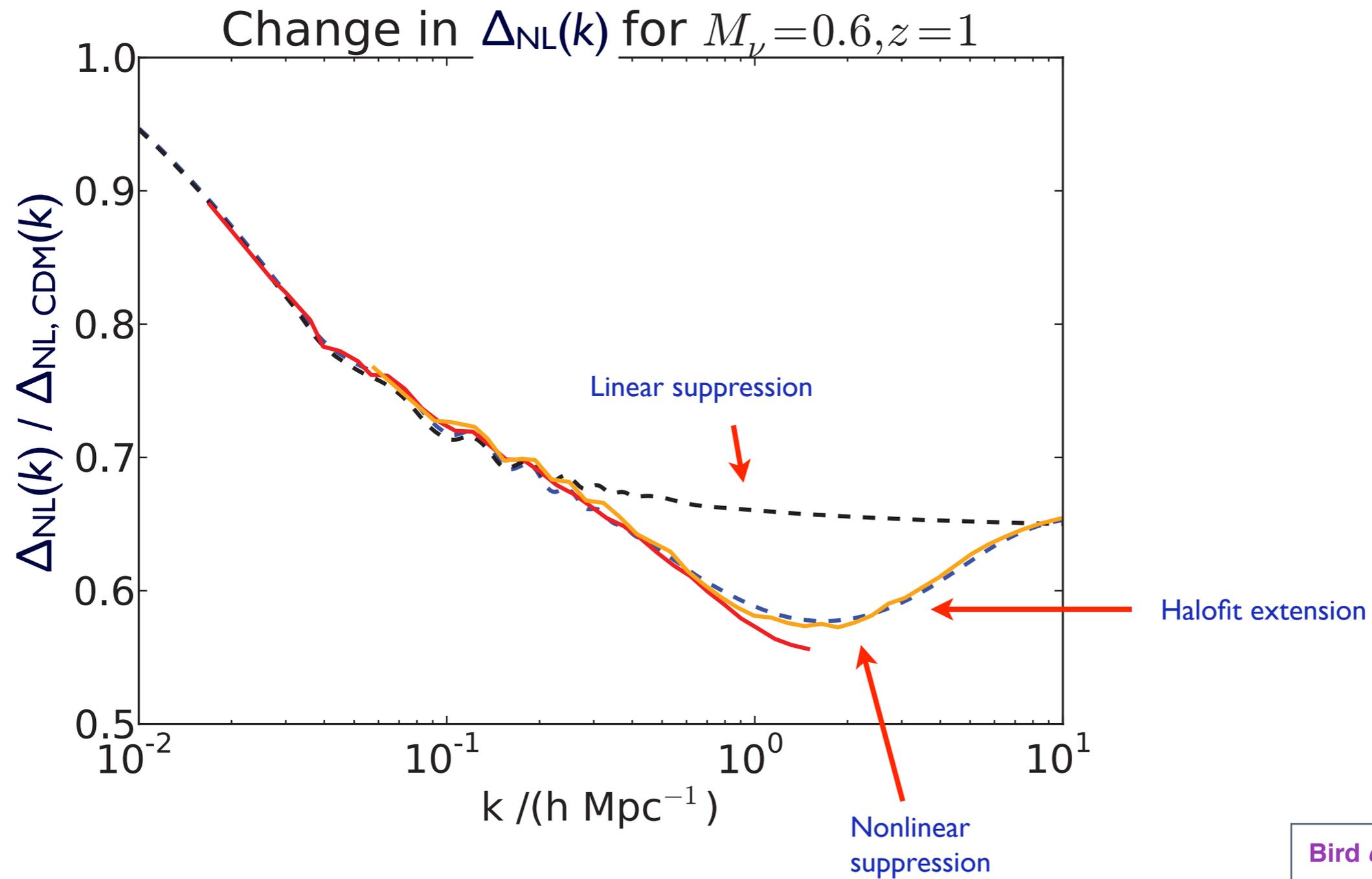
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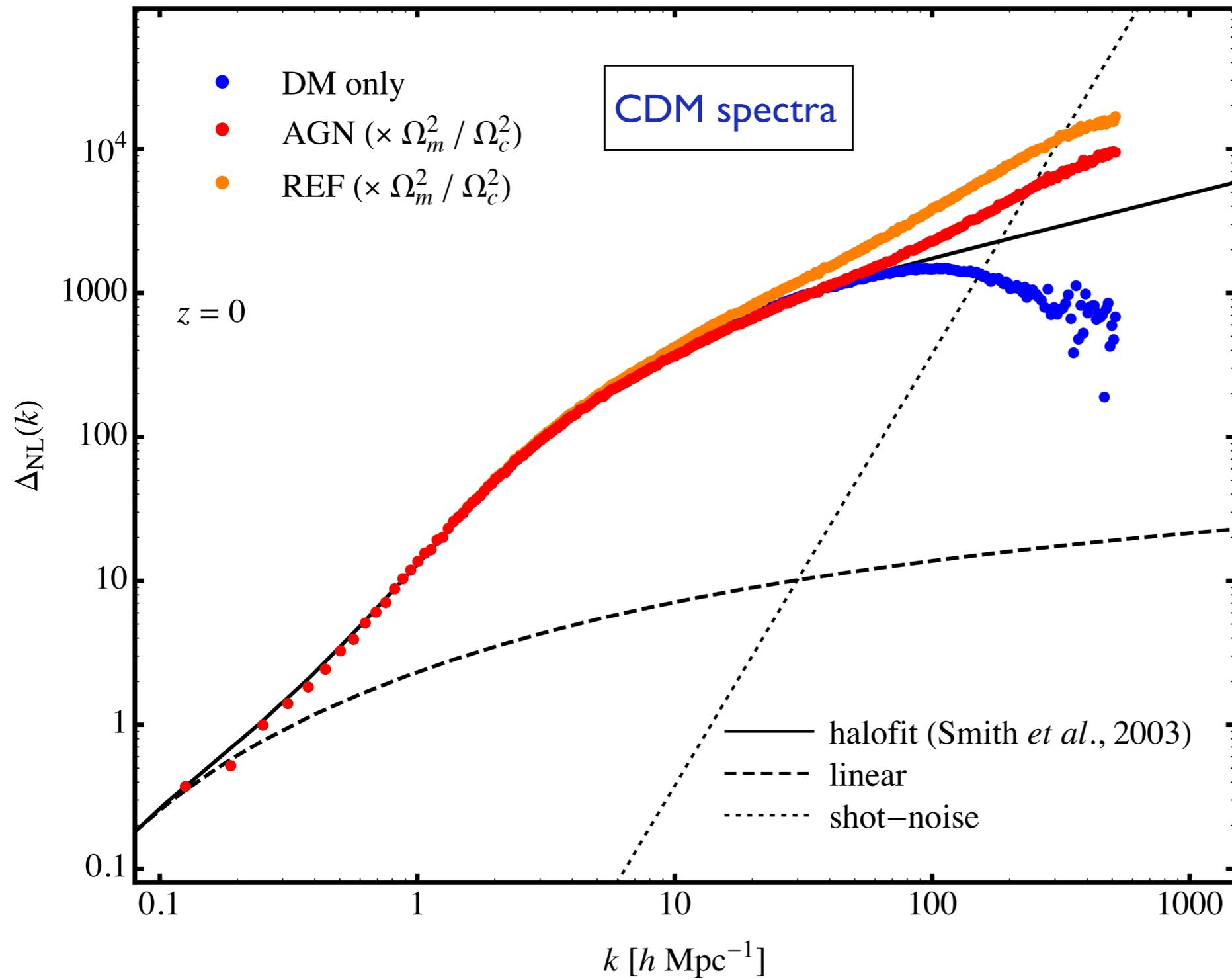
Estimating uncertainties ...



Other uncertainties: neutrinos



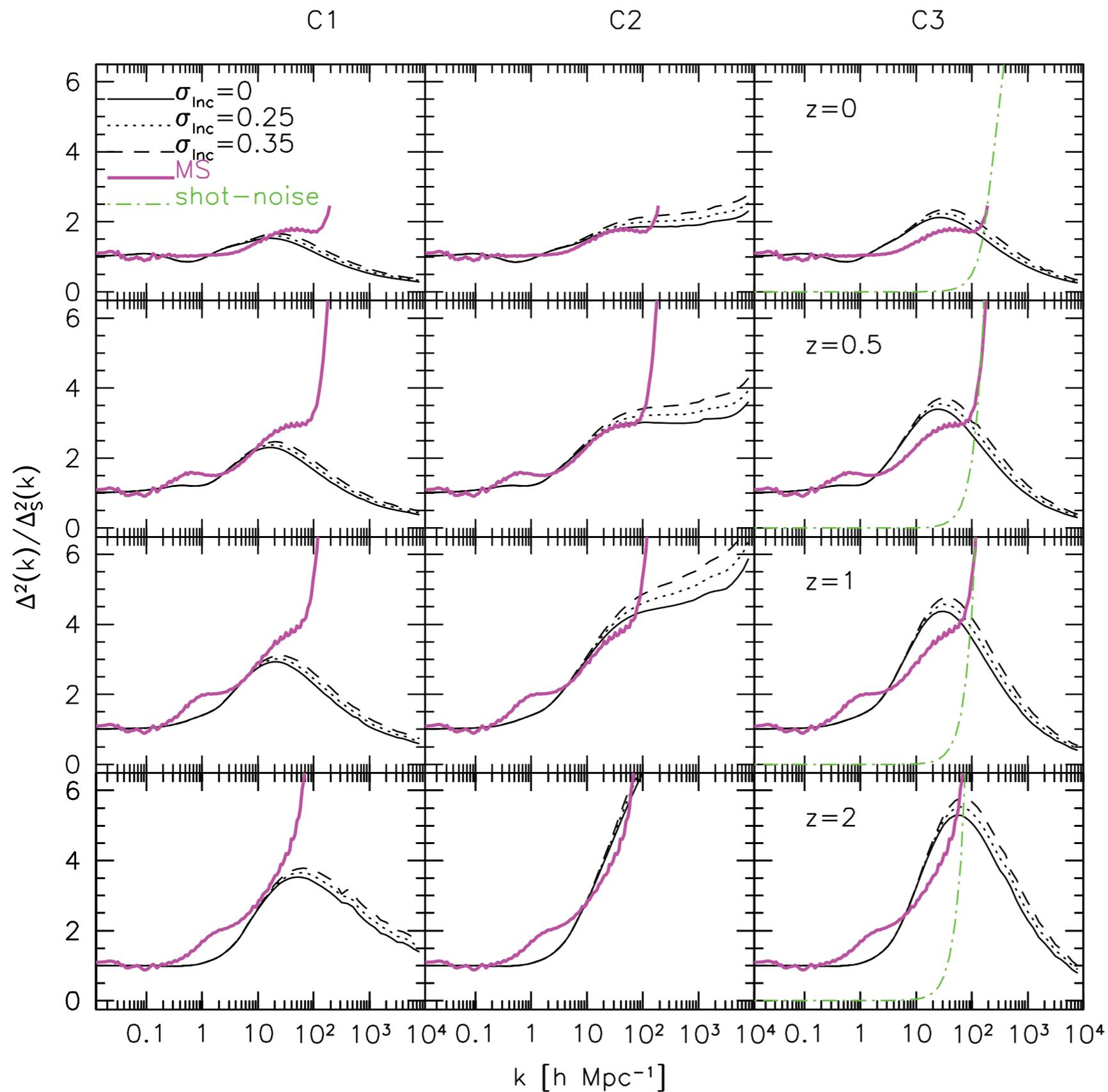
Other uncertainties: baryons



data from

van Daalen *et al.* (2011)

The nonlinear Power Spectrum *from* the Halo Model



Conclusion

Halo Model:

$$\zeta(z) = \frac{1}{\Omega_m \rho_c} \int_{M_{min}} dM \frac{dn}{dM} M \frac{\Delta_{vir}(z)}{3} \langle F \rangle \quad F = c_v^3(M, z) \frac{\int_0^{c_v} dx x^2 \kappa^2(x)}{[\int_0^{c_v} dx x^2 \kappa(x)]^2}$$

+ subhalos properties (x2)

or

Power Spectrum:

$$\zeta(z) = \langle \delta^2(z, \hat{\Omega}) \rangle = \int_0^{k_{max}} \frac{dk}{k} \frac{k^3 P_{NL}(k)}{2\pi^2}$$

The DM extragalactic annihilation flux can be computed in the Halo Model from 3 or more quantities determined from simulations or directly from the Power Spectrum, with minimal assumptions